

The E&M Equations, Explained

Every core relationship with what it means and exactly when to reach for it

Knowing an equation is not knowing when to use it. This guide walks the E&M toolkit in the order the physics builds, from a single charge to Maxwell's induction, with an original plain-English note on each. Read it as a map, not a list.

■ Part 1. Electrostatics

fields, flux, potential

Force, field and Gauss

$$\bullet F = \frac{kq_1q_2}{r^2}, \quad \mathbf{E} = \frac{\mathbf{F}}{q} = \frac{kq}{r^2}\hat{\mathbf{r}}$$

Coulomb's law and the field of a point charge. Use for point charges or as the building block for an integral over a continuous distribution.

$$\bullet \oint \mathbf{E} \cdot d\mathbf{A} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

Gauss's law. Use *instead of* integrating when the charge has spherical, cylindrical, or planar symmetry, it makes the field fall out in one line.

Potential

$$\bullet V = \frac{kq}{r}, \quad U = qV, \quad \mathbf{E} = -\frac{dV}{dr}, \quad \Delta V = -\int \mathbf{E} \cdot d\mathbf{l}$$

Potential is energy per charge. Go from V to E by differentiating, from E to V by integrating. Potential is a *scalar*, so it adds without vector components, often easier than fields.

■ Part 2. Capacitance and circuits

storing charge and driving current

Capacitors

$$\bullet C = \frac{Q}{V}, \quad C = \frac{\epsilon_0 A}{d}, \quad U = \frac{1}{2}CV^2 = \frac{Q^2}{2C}$$

Capacitance depends on geometry only. Series capacitors add reciprocally ($\frac{1}{C} = \sum \frac{1}{C_i}$); parallel add directly, the opposite of resistors.

Current, resistance, power

$$\bullet I = \frac{dQ}{dt}, \quad R = \frac{\rho L}{A}, \quad V = IR, \quad P = IV = I^2R = \frac{V^2}{R}$$

Ohm's law and power. Series resistors add; parallel add reciprocally. Choose the power form that uses the two quantities you already know.

RC circuits

- $q(t) = Q(1 - e^{-t/RC})$ (charging), $q(t) = Q e^{-t/RC}$ (discharging), $\tau = RC$

The time constant $\tau = RC$ is the time to reach $\approx 63\%$ of the change. At $t = 0$ a capacitor acts like a wire; after a long time, like a break in the circuit.

■ Part 3. Magnetism and induction

moving charge and changing flux

Magnetic force and Ampère

- $\mathbf{F} = q\mathbf{v} \times \mathbf{B}$, $\mathbf{F} = I\mathbf{L} \times \mathbf{B}$

Force on a moving charge and on a current. The cross product means the force is *perpendicular* to both $\mathbf{v}$ and $\mathbf{B}$; use the right-hand rule and note it does no work.

- $\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I_{\text{enc}}$, $B_{\text{solenoid}} = \mu_0 n I$

Ampère's law, the magnetic analogue of Gauss. Use it for the field of a long wire, a solenoid, or a toroid, wherever the current has symmetry.

Faraday's law and inductance

- $\Phi_B = \int \mathbf{B} \cdot d\mathbf{A}$, $\varepsilon = -\frac{d\Phi_B}{dt}$

Faraday's law: a *changing* magnetic flux drives an EMF. The minus sign (Lenz's law) says the induced current opposes the change that made it.

- $\varepsilon = -L \frac{dI}{dt}$, $U = \frac{1}{2} L I^2$

An inductor opposes changes in current and stores energy in its field. In an RL circuit the current rises with time constant $\tau = L/R$.

The through-line. Gauss and Ampère both trade a hard integral for symmetry. Potential (a scalar) is often the easy road to a field. And induction is the one place a *rate of change* drives the physics, watch for “changing” in the question.

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